

Solutions to Chen's Plasma Physics

Kalman Krizhnik

1-1. Compute the density (in units of m^{-3}) of an ideal gas under the following conditions:

a) At 0°C and 760 Torr pressure (1 Torr = 1mm Hg). This is called the Loschmidt number.

b) In a vacuum at 10^{-3} Torr at room temperature (20°C). This number is a useful one for the experimentalist to know by heart (10^{-3} Torr = 1 micron).

a) Avogadro's number is $N_A = 6.022 \times 10^{23}$. One mole of gas at STP occupies 22.4 liters. 1 liter is 1×10^{-3} cubic meters. Thus, the number per cubic meter is $N_A/n = 6.022 \times 10^{23}/(22.4 \times 10^{-3}) = 2.66 \times 10^{25} m^{-3}$. Thus, the Loschmidt number is 2.66×10^{25} $\square$

b) Using $PV=NkT$, we obtain (with $R = 1.4 \times 10^{-23} \text{ J } K^{-1}$ and 1 Torr = 133 Pa):

$$n = \frac{N}{V} = \frac{P}{kT} = \frac{10^{-3} \times 133}{1.4 \times 10^{-23} \times (20 + 273)} = 3.3 \times 10^{19} m^{-3} \quad \square \quad (1)$$

1-2. Derive the constant A for a normalized one-dimensional Maxwellian distribution

$$f(u) = Ae^{-mu^2/2kT} \quad (2)$$

such that

$$\int_{-\infty}^{\infty} f(u) du = 1 \quad (3)$$

This one is straightforward. Just integrate:

$$1 = \int_{-\infty}^{\infty} Ae^{-mu^2/2kT} du = A \sqrt{\frac{2\pi kT}{m}} \Rightarrow A = \sqrt{\frac{m}{2\pi kT}} \quad \square \quad (4)$$

1-4. Compute the pressure, in atmospheres and in tons/ft^2 , exerted by a thermonuclear plasma on its container. Assume $kT_e = kT_i = 20\text{keV}$, $n = 10^{21} m^{-3}$ and $p = nkT$, where $T = T_i + T_e$.

This is just unit conversion, albeit with units that nobody really ever remembers. For reference, $1 \text{ keV} = 1.6 \times 10^{-19} \text{ J}$, so

$$p = 10^{21} \times (20\text{keV} + 20\text{keV}) = 4 \times 10^{22} m^{-3} \text{ keV} = 4 \times 10^{22} m^{-3} \text{ J} = 4 \times 10^3 \text{ N}/m^2 \quad (5)$$

But $1 \text{ atm} = 10^5 \text{ N}/m^2 = 1 \text{ ton}/ft^2$, so (Note: I think there is a mistake in Chen's solutions here. If I am mistaken, please let me know).

$$p = 0.04 \text{ atm} = 0.04 \text{ ton}/ft^2 \quad \square \quad (6)$$

1-5. In a strictly steady state situation, both the ions and the electrons will follow the Boltzmann relation

$$n_i = n_0 e^{-e\phi/kT_i} \quad (7)$$

For the case of an infinite, transparent grid charged to a potential ϕ , show that the shielding distance is given approximately by

$$\lambda_D^{-2} = \frac{ne^2}{\epsilon_0} \left(\frac{1}{kT_e} + \frac{1}{kT_i} \right) \quad (8)$$

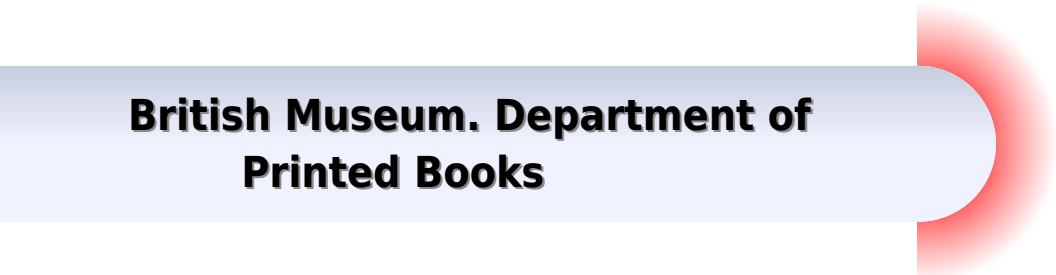
Show that λ_D is determined by the temperature of the colder species.

We'll use Poisson's equation

$$\nabla^2 \phi = \frac{en_e - en_i}{\epsilon_0} = \frac{en_0}{\epsilon_0} (e^{e\phi/kT_e} - e^{-e\phi/kT_i}) \approx \frac{en_0}{\epsilon_0} \left(1 + \frac{e\phi}{kT_e} - 1 - \frac{-e\phi}{kT_i} \right) = \frac{e^2 n}{\epsilon_0} \left(\frac{\phi}{kT_e} + \frac{\phi}{kT_i} \right) \quad (9)$$

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